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**The Crystal Structure of Pyridinium Tetracosabromoantimon(III)triantimon(V)ate,  $(C_5H_5NH)_6Sb^{III}Sb^V_3Br_{24}$ <sup>1</sup>**BY STEPHEN L. LAWTON,<sup>\*2</sup> ROBERT A. JACOBSON,<sup>3</sup> AND ROBERT S. FRYE<sup>4</sup>

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Pyridinium tetracosabromoantimon(III)triantimon(V)ate,  $(C_5H_5NH)_6Sb^{III}Sb^V_3Br_{24}$ , crystallizes as intensely deep red acicular crystals in the orthorhombic space group  $Cmcm$  ( $D_{2h}^{17}$ ) with unit cell parameters  $a = 17.474 \pm 0.004$ ,  $b = 23.697 \pm 0.006$ , and  $c = 16.390 \pm 0.006$  Å. The observed and calculated densities are  $2.7 \pm 0.2$  and  $2.824 \pm 0.002$  g/cm<sup>3</sup>, respectively. The structure was refined by least-squares methods to a final conventional  $R$  index of 0.074 using three-dimensional X-ray diffraction counter data. Three crystallographically independent pyridinium cations, one  $Sb^{III}Br_6^{3-}$ , and two independent  $Sb^VBr_6^-$  anions form an asymmetric unit. All three anions exhibit slight deviations from octahedral ( $O_h$ ) symmetry. One of the  $Sb^VBr_6^-$  ions forms nearly linear  $\cdots Br-Sb^V-Br \cdots Br-Sb^V-Br \cdots$  chains extending along the  $c$  direction of the unit cell. The other two ions,  $Sb^VBr_6^-$  and  $Sb^{III}Br_6^{3-}$ , form nearly linear  $\cdots Br-Sb^V-Br \cdots Br-Sb^{III}-Br \cdots$  chains extending along the  $a$  direction of the unit cell. The two chains are nonintersecting. Bromine $\cdots$ bromine contacts along both chains are quite short, being  $3.243 \pm 0.007$  Å within the former chain and  $3.486 \pm 0.007$  Å within the latter; both are considerably shorter than the anticipated 3.90-Å van der Waals radius sum. All cations are oriented such that the plane of each ring is nearly parallel to the unit cell  $c$  axis and essentially normal to one or more of the many  $Br \cdots Br$  contacts. The derived  $Sb^V-Br$  and  $Sb^{III}-Br$  bond lengths, corrected for thermal effects assuming rigid-body libration of the ions, average  $2.546 \pm 0.011$  and  $2.799 \pm 0.007$  Å, respectively. A brief comparison of these bond lengths is made with other related structures.

**Introduction**

This paper describes another in a series of  $R_xSb_yX_z$  ( $R =$  cation;  $X = Cl, Br, I$ ) structures. Those previously reported include  $(NH_4)_4Sb^{III}Sb^VBr_{12}$ ,<sup>5,6</sup>  $Rb_4Sb^{III}Sb^VBr_{12}$ ,<sup>7</sup>  $(C_6H_7NH)_2Sb^VBr_9$ ,<sup>8</sup>  $(C_5H_5NH)_6Sb^{III}Sb^VBr_{24}$ ,<sup>9</sup> and  $(C_5H_5NH)Sb^{III}Cl_4$ .<sup>10</sup> We now wish to report the results of a detailed single-crystal structure determination of another pyridinium analog, one containing Sb(III) and Sb(V),  $(C_5H_5NH)_6Sb^{III}Sb^V_3Br_{24}$ . This material is intensely colored, appearing deep red to black, within the full temperature range from its melting point (201–202°) down to the temperature of liquid nitrogen (–173°). This coloration is unlike that of the 2-, 3-, and 4-methylpyridinium derivatives,  $(C_6H_7NH)_2SbBr_9$ , which are similarly intensely colored at their melting points (range 116–38°) and room temperature but are bright orange to red at –173°.<sup>11</sup> It was the purpose of the present investigation to compare this structure with that of the 2-methylpyridinium derivative previously reported<sup>8</sup> and to relate differences in structure with differences in their charge-transfer absorption. These antimony halide complexes afford good model systems for the study of weak interactions in the solid state; also the geometry of possible Sb(III) moieties is

of considerable interest, especially in relation to the effect of the lone pair. The study augments two previous single crystal structure investigations of the Sb(III)–Sb(V) systems,<sup>5–7</sup> as well as three structure determinations of ionic antimony halides with  $C_5H_5NH^+$  as a cation.<sup>9,10,12</sup> Physical and chemical properties of this material and other intensely colored  $R_xSb_yBr_z$  complexes synthesized with quaternary ammonium cations and saturated and unsaturated heterocyclic amines have already been discussed.<sup>5,11</sup>

**Experimental Section**

**Preparation.**—The salt was prepared by treating 1.0 ml of pyridine with a solution of 2.50 g of antimony tribromide and 0.5 ml of liquid bromine in 10 ml of warm concentrated hydrobromic acid (48%). The black crystalline precipitate was filtered by vacuum filtration through a sintered-glass funnel, rinsed with a small amount of cold, concentrated hydrobromic acid (48%), and dried on a porous porcelain plate in a desiccator containing concentrated sulfuric acid as the desiccant and a small partial pressure of bromine vapor.

Single crystals suitable for the X-ray investigation were prepared by recrystallizing the raw material from 2–3 ml of hot, concentrated hydrobromic acid (48%) containing 1–2 drops of concentrated sulfuric acid and 1–2 drops of liquid bromine, following a procedure described elsewhere.<sup>8,11</sup>

**Anal.** Calcd for  $(C_5H_5NH)_6Sb^{III}Sb^V_3Br_{24}$ : C, 13.22; N, 3.08; H, 1.33; Sb, 17.87; Br, 64.50. Calcd for  $(C_5H_5NH)_3Sb_2Br_{12}$ : C, 12.49; N, 2.91; H, 1.26; Sb, 16.88; Br, 66.46. Found: C, 13.17; N, . . . ; H, 1.64; Sb, 16.64; Br, 63.7.

**Crystal Data.**—Pyridinium tetracosabromoantimon(III)triantimon(V)ate,  $(C_5H_5NH)_6Sb^{III}Sb^V_3Br_{24}$  (formula weight 2885.48) crystallizes in the orthorhombic space group  $Cmcm$  ( $D_{2h}^{17}$ ); lattice parameters at 23° are  $a = 17.474 \pm 0.004$ ,  $b = 23.697 \pm 0.006$ , and  $c = 16.390 \pm 0.006$  Å;  $V = 6787 \pm 6$  Å<sup>3</sup> ( $23 \pm 3^\circ$ );  $d_{\text{obsd}} = 2.7 \pm 0.2$  g/cm<sup>3</sup> (by flotation) and  $d_{\text{calcd}} = 2.824 \pm 0.002$  g/cm<sup>3</sup> for  $Z = 4$  formula units of  $(C_5H_5NH)_6Sb_4Br_{24}$  per unit cell. The color is jet black (reflected light) to deep red (transmitted light from very thin sections of the crystal); the crystal habit is acicular, elongated along the  $c$  direction.

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(1) Work was performed in part at the Ames Laboratory of the U. S. Atomic Energy Commission and Mobil Research and Development Corp. Ames Laboratory Contribution No. 2172.

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(10) (a) S. K. Porter and R. A. Jacobson, *Chem. Commun.*, 1244 (1967);

(b) S. K. Porter and R. A. Jacobson, *J. Chem. Soc. A*, 1356 (1970).

(11) M. L. Hackert, S. L. Lawton, and R. A. Jacobson, *Proc. Iowa Acad. Sci.*, **75**, 97 (1968).

The crystal symmetry was determined from Weissenberg and precession photographs which yielded systematic extinctions ( $hkl$ ,  $h + k = 2n + 1$ , and  $h0l$ ,  $l = 2n + 1$ ) consistent with the space groups  $Cmc2_1$  ( $C_{2v}^{12}$ ),  $C2cm$  ( $C_{2v}^{16}$ ), and  $Cmcm$  ( $D_{2h}^{17}$ ). Intensity statistics strongly favored the centrosymmetric space group as the correct choice,<sup>13</sup> later confirmed by the successful refinement of the derived structure. The lattice parameters were determined by a least-squares fit<sup>14</sup> to weighted high-angle copper  $K\alpha_1$  ( $\lambda$  1.54050 Å) and  $K\alpha_2$  ( $\lambda$  1.54434 Å) reflections measured from  $hk0$  and  $0kl$  Weissenberg zones calibrated with superimposed aluminum powder lines ( $a_0 = 4.0330$  Å). The Nelson-Riley extrapolation function was employed in the refinement. Indicated errors in the cell parameters are  $2\sigma$ .

**Collection and Reduction of X-Ray Intensity Data.**—Complete three-dimensional X-ray diffraction intensity data were taken at room temperature with zirconium-filtered molybdenum  $K\alpha$  radiation from a crystal 0.373 mm long and of cross section  $0.076 \times 0.138$  mm ( $c^*$ ,  $a^*$ , and  $b^*$ , respectively). The crystal was mounted in a 0.3-mm thin-walled Lindemann glass capillary with the needle axis (unit cell  $c$  axis) coincident with the  $\phi$  axis of the diffractometer. A Siemens diffractometer equipped with a scintillation detector and pulse height discriminator was used with the moving-crystal, moving-counter measurement technique ( $\theta$ ,  $2\theta$  coupling) and a  $3.5^\circ$  takeoff angle. The receiving aperture size selected to minimize extraneous background was 5.0 mm wide by 5.0 mm high. The counter angle,  $2\theta$ , was scanned over  $2^\circ$  at a speed of  $2^\circ/\text{min}$ . Background counts of 12 sec were taken at each end of the  $2\theta$  scan. All scans were recorded on a chart recorder to provide visual evidence for the existence of observed reflections, proper peak shape, reflection centering in  $2\theta$ , and nonoverlap of adjacent reflections. A total of 3255 independent reflections were measured within the range  $2\theta \leq 50^\circ$ . Equivalent reflections were not measured. Typical background counts at 10, 20, 30, 40, and  $50^\circ 2\theta$  were 37, 8, 4, 3, and 2 counts/sec, respectively. Three standard reflections were measured periodically as a check on electronic and crystal stability and a 5% increase in intensities was observed throughout the 60-hr recording period. The alignment of the crystal was checked on a daily basis for  $\phi$  independence at  $\chi = 90^\circ$  and adjusted when necessary.

The mosaicity of the crystal was examined by means of a narrow-source (takeoff angle  $0.5^\circ$ )  $2\theta$ -scan technique at  $2\theta < 11^\circ$ . In this region, the  $2\theta$ -scan and  $\omega$ -scan techniques yield comparable results.<sup>15</sup> Widths at half-maximum for three typical strong noncoplanar reflections ranged from 0.09 to  $0.16^\circ \theta$ .

The raw intensity of each reflection was corrected for background, Lorentz, polarization, and absorption effects and for the slight increase in intensities as monitored by the reference reflections. Transmission factors were calculated by the program ACACA.<sup>16</sup> Because of the large linear absorption coefficient ( $\mu = 164.3 \text{ cm}^{-1}$ ) and the variation in  $\mu R$  (0.7–1.2) with the variation in the needle radius  $R$ , refined<sup>17</sup> crystal dimensions normal to the needle axis, rather than those obtained by optical measurement, were used in the absorption correction. Pursuant to this refine-

ment, the shape, or profile, of the needle cross section was first carefully examined by measuring the intensity of a strong low-order  $00l$  reflection, in this case  $004$ , at  $\chi = 90^\circ$  in  $15^\circ$  intervals of  $\phi$  from 0 to  $360^\circ$ . These intensities were used to prepare a plot of  $I_n/I_{av}$  vs.  $\phi$ , where  $I_n$  (corrected for background) was the intensity at  $\phi = n$  and  $I_{av}$  was the average intensity over the full range of  $\phi$ . By systematically adjusting the shape and optically measured crystal thickness, a "best fit" of the transmission factor curve  $A_n/A_{av}$  with that of  $I_n/I_{av}$  plotted against  $\phi$  was obtained by defining the crystal as an ellipsoidal cylinder with *minor* diameter 0.0844 mm and *major* diameter 0.1494 mm. The maximum deviation in  $F_o$  from the mean, defined as  $100(|F_{\text{max}} - F_{\text{min}}|/F_{av})$ , for the 24 independent measurements of the  $004$  reflection was subsequently reduced from 29.8% with no applied correction to only 1.2% with this correction. The new cross section was further tested on the  $008$  reflection for 24 values of  $\phi$  and found to be satisfactory. The calculated ratio of the minimum and maximum transmission factors for the full set of three-dimensional data was 0.397. Effects of secondary extinction proved not to be a major problem and so no such correction was applied.

The estimated error in each intensity measurement was calculated by the expression<sup>18</sup>  $\sigma(I) = [C_T + 0.25(t_c/t_b)^2(B_1 + B_2) + (pK)^2]^{1/2}$  where  $C_T$  is the total integrated peak count obtained in a scan time  $t_c$ ,  $B_1$  and  $B_2$  are the background counts each obtained in time  $t_b$ , and  $K = C_T - 0.5(t_c/t_b)(B_1 + B_2)$ . The value of  $p$  was selected as 0.06. Each  $\sigma(I)$  was then corrected for Lorentz, polarization, and absorption effects. The estimated standard deviation in each  $F_o$  was calculated by the expression  $\sigma(F_o) = [(I + \sigma(I))^2 - |F_o|^2]^{1/2}$ , a function based on the finite difference method.<sup>19</sup> These standard deviations were used during the least-squares refinements to weight the observed structure factors. A total of 1421 reflections were observed above the background level of which 1173 had  $F_o^2 \geq \sigma(F_o^2)$ . Those with  $F_o^2 < \sigma(F_o^2)$  were considered as unobserved and thus omitted from the refinements.

### Solution and Refinement of the Structure

Preparation of this salt was first reported by Petzold in 1933.<sup>20</sup> The empirical formula was given as  $R_3\text{Sb}_2\text{Br}_{12}$  ( $R = \text{pyridinium}$ ). Analysis of our material indicated two plausible formulas,  $R_3\text{Sb}_2\text{Br}_{11}$  and  $R_3\text{Sb}_2\text{Br}_{12}$ , with calculated densities, based on 8 formula units/unit cell, of  $2.667 \pm 0.002$  and  $2.824 \pm 0.002 \text{ g/cm}^3$ , respectively. Our observed density was  $2.7 \pm 0.2 \text{ g/cm}^3$ .

The data were sharpened using the method of Jacobson, *et al.*,<sup>21</sup> and a sharpened three-dimensional Patterson function was computed.<sup>22</sup> Analysis of the Patterson-Harker sections and subsequent Fourier syntheses revealed three crystallographically independent  $\text{SbBr}_6$  octahedra compatible with space group  $Cmc2_1$  and  $Cmcm$ . Bond length calculations indicated the presence of trivalent and pentavalent antimony in the ratio 1:3. Careful analysis of the trivalent anion through refinement procedures in the space group  $Cmc2_1$  ruled out disordered  $\text{Sb}^{\text{III}}\text{Br}_5^{2-}$  in which the lone pair

(13) (a) H. Lipson and W. Cochran, "The Determination of Crystal Structures," G. Bell and Sons, London, 1957, pp 32–41; (b) L. Guggenberger, "WSTAT, A Fortran Crystallographic Intensity Statistical Analysis Program," Experimental Station, E. I. du Pont de Nemours and Co., Wilmington, Del., 1967.

(14) D. E. Williams, "LCR-2, A Fortran Lattice Constant Refinement Program," USAEC Report IS-1052, Ames Laboratory, Iowa State University, Ames, Iowa, 1964.

(15) T. C. Furnas, "Single Crystal Orienter Instruction Manual," General Electric Co., Milwaukee, Wis., 1966.

(16) (a) B. J. Wuensch and C. T. Prewitt, *Z. Kristallogr., Kristallgeometrie, Kristallphys., Kristallchem.*, **122**, 24 (1965); (b) C. T. Prewitt, "ACACA, A Fortran Polyhedral Absorption Correction Program," Experimental Station, E. I. du Pont de Nemours and Co., Wilmington, Del., 1966.

(17) Initial dimensions of the air-sensitive crystal were obtained by optical measurement with a  $100\times$  microscope while the crystal was in the capillary and therefore tend to be less reliable than those obtained without the interfering refraction problems arising from the capillary walls and crystal adherent. The method described here was devised to get around these problems. It has been successfully applied to a number of crystals and found to result in meaningful improvements in the values of the transmission factors.

(18) P. W. R. Corfield, R. J. Doedens, and J. A. Ibers, *Inorg. Chem.*, **6**, 197 (1967).

(19) D. E. Williams and R. E. Rundle, *J. Amer. Chem. Soc.*, **86**, 1660 (1964).

(20) W. Petzold, *Z. Anorg. Allg. Chem.*, **215**, 92 (1933).

(21) R. A. Jacobson, J. A. Wunderlich, and W. N. Lipscomb, *Acta Crystallogr.*, **14**, 598 (1961).

(22) In addition to various local programs for the IBM 7074 and CDC 1604 computers, programs used in the solution and refinement of this structure were Fitzwater, Benson, and Jackobs' Fourier program, Guggenberger's FOUR Fourier program, Ibers and Doedens' NUCLS crystallographic least-squares group-refinement program, Busing and Levy's ORFFE function and error program, and Johnson's ORTEP thermal ellipsoid plotting program.

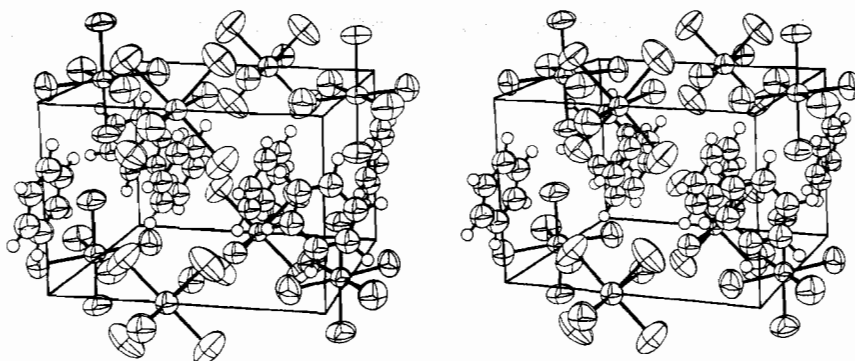


Figure 1.—A stereographic view of the packing of ions in crystals of  $(C_5H_5NH)_6Sb^{III}Sb^V_3Br_{24}$ . The view depicts the contents of one-eighth the unit cell (two asymmetric units). Ellipsoidal boundaries of the atoms (except carbon and hydrogen) are at the 80% probability level.

of electrons might occupy at random the four possible octahedral sites on a mirror plane. The structure was therefore assumed to be  $6C_5H_5NH^+ \cdot Sb^{III}Br_6^{3-} \cdot 3Sb^V \cdot Br_6^-$ . A series of additional three-dimensional electron density syntheses ( $F_o$  and  $F_o - F_c$ ) and isotropic least-squares refinements led to the location of all cations. A final difference map revealed that all electron density had now been accounted for and that the empirical formula,  $(C_5H_5NH)_6Sb_4Br_{24}$ , had been established, confirming Petzold's original analysis.

The structure was refined by using a full-matrix, least-squares procedure. The function minimized was  $\sum w(|F_o| - |F_c|)^2$ , where  $|F_o|$  and  $|F_c|$  are the observed and calculated structure amplitudes, respectively, and  $w$  is the weight defined as  $1/\sigma^2(F_o)$ . Atomic scattering factors for neutral atoms tabulated by Hanson, *et al.*,<sup>23</sup> were used. Anomalous parts of the Sb and Br scattering factors were obtained from Templeton's tabulation<sup>24,25</sup> and were included in the calculated structure factors.<sup>26</sup> All initial Fourier syntheses and least-squares refinements were carried out in the noncentrosymmetric space group  $Cmc2_1$  to avoid any assumptions about extra mirror symmetry, but after all constituent ions had been identified and trial least-squares results evaluated, it became apparent that the more probable space group was  $Cmcm$ . A subsequent statistical analysis of the intensities<sup>13</sup> further indicated the presence of a center of symmetry.

Initially a fully isotropic refinement in  $Cmcm$  was carried out with the pyridinium rings constrained to  $D_{6h}$  symmetry using the least-squares group refinement procedure described by La Placa and Ibers.<sup>27</sup> In this refinement C-C and C-H distances of 1.376 and 0.98 Å, respectively, were assumed. Owing to the cavity-like environment in which the cations were situated, ordered nitrogen positions were considered highly unlikely and so no such ordering was assumed; all ring

(23) H. P. Hanson, F. Herman, J. D. Lea, and S. Skillman, *Acta Crystallogr.*, **17**, 1040 (1964).

(24) D. H. Templeton, "International Tables for X-Ray Crystallography," Vol. III, Kynoch Press, Birmingham, England, 1962, pp 215, 216, Table 3.3.2C.

(25) For Mo K $\alpha$  radiation the real and imaginary dispersion corrections,  $\Delta f'$  and  $\Delta f''$ , for atomic scattering factors are -0.6 and 2.0 for antimony and -0.3 and 2.6 for bromine, both at a  $(\sin \theta)/\lambda$  value of 0.0.<sup>24</sup> The effects on the atomic scattering factors caused by these values are considered to be moderately significant.

(26) J. A. Ibers and W. C. Hamilton, *Acta Crystallogr.*, **17**, 781 (1964).

(27) S. J. La Placa and J. A. Ibers, *ibid.*, **18**, 511 (1965).

atoms were treated as carbon. The C-C bond length was taken as a weighted average of two C-N lengths, 1.340 Å each,<sup>28</sup> and four C-C lengths, 1.394 Å each, assuming nitrogen statistically occupied one-sixth of each ring position. It is now fairly well established that carbon-hydrogen distances obtained by X-ray diffraction techniques are systematically shorter than those determined from spectroscopic, electron diffraction, or neutron diffraction studies.<sup>29,30</sup> Accordingly, the C-H bond length of 0.98 Å, rather than the spectroscopically determined value of 1.084 (5) Å for an aromatic carbon-hydrogen bond,<sup>31</sup> was used.<sup>32</sup> Each ring was assigned a single, variable isotropic thermal parameter and six variable positional parameters.<sup>27</sup> The origin of the internal system was taken at the ring center with  $a_3'$  normal to the ring and  $a_1'$  intersecting a vertex. This initial refinement of the rings together with the other atoms, each of which was assigned a variable isotropic thermal parameter, converged rapidly to  $R$  values of  $R_1 = \sum (|F_o| - |F_c|) / \sum |F_o| = 0.129$  and  $R_2 = (\sum w(|F_o| - |F_c|)^2 / \sum w|F_o|^2)^{1/2} = 0.075$ .

A difference map based on the preceding isotropic refinement provided evidence for anisotropic thermal motion of the heavy atoms which seemed physically reasonable, and ellipsoidal thermal parameters were introduced for antimony and bromine. Analysis of mean  $w\Delta^2$  ( $\Delta = |F_o| - |F_c|$ ) as functions of  $|F_o|$  and  $(\sin \theta)/\lambda$  at this point revealed a trend in which the weak reflections were being slightly underweighted and the weights were subsequently modified to remove this dependency.

Convergence was reached with  $R_1 = 0.074$  and  $R_2 = 0.061$ . The corresponding values for all 1421 reflections were  $R_1 = 0.100$  and  $R_2 = 0.078$ . The final standard deviation for an observation of unit weight (*i.e.*, the "error of fit") was 1.20, where the "error of fit" is defined by  $[\sum w(|F_o|^2 - |F_c|^2) / (n - m)]^{1/2}$  with  $n$  being the number of observations (1173) and  $m$  the number of

(28) B. Bak, L. Hansen-Nygaard, and J. Rastrup-Anderson, *J. Mol. Spectrosc.*, **2**, 361 (1958).

(29) W. C. Hamilton and J. A. Ibers, "Hydrogen Bonding in Solids: Methods of Molecular Structure Determinations," W. A. Benjamin, New York, N. Y., 1968.

(30) B. Dawson, *Aust. J. Chem.*, **18**, 595 (1965).

(31) L. E. Sutton, Ed., *Chem. Soc., Spec. Publ.*, No. 18, S18s (1965).

(32) See, for example: (a) V. R. Magnuson and G. D. Stucky, *Inorg. Chem.*, **8**, 1427 (1969); (b) M. R. Churchill and F. R. Scholer, *ibid.*, **8**, 1950 (1969).

TABLE I  
 FINAL POSITIONAL, THERMAL, AND GROUP PARAMETERS FOR  $(C_5H_5NH)_6Sb_4Br_{24}$ 

Atom	Position	$x^a$	$y$	$z$	$\beta_{11}^b$	$\beta_{22}$	$\beta_{33}$	$\beta_{12}$	$\beta_{13}$	$\beta_{23}$	Iso equiv $B,^c \text{ \AA}^2$
Sb(1)	8g	0.22100 (20)	0.01773 (15)	0.25* <sup>d</sup>	300 (13)	157 (6)	245 (12)	-1 (10)	0*	0*	3.27
Sb(2)	4c	0.0*	0.28391 (20)	0.25*	179 (19)	149 (10)	387 (25)	0*	0*	0*	3.23
Sb(3)	4c	0.0*	0.25334 (10)	0.75*	179 (17)	128 (10)	300 (22)	0*	0*	0*	2.76
Br(1)	8g	0.11683 (38)	0.09191 (26)	0.25*	412 (30)	286 (16)	571 (36)	133 (18)	0*	0*	5.86
Br(2)	8g	0.12443 (39)	-0.06194 (25)	0.25*	399 (29)	267 (15)	579 (35)	-83 (17)	0*	0*	5.70
Br(3)	8g	0.32997 (37)	-0.05483 (24)	0.25*	384 (27)	220 (14)	558 (34)	94 (16)	0*	0*	5.21
Br(4)	8g	0.32143 (41)	0.09421 (25)	0.25*	542 (33)	235 (15)	563 (37)	-118 (18)	0*	0*	5.99
Br(5)	16h	0.22121 (25)	0.01785 (18)	0.40448 (20)	517 (18)	319 (9)	258 (14)	10 (14)	-28 (14)	17 (12)	5.42
Br(6)	8g	0.14504 (30)	0.28170 (23)	0.25*	180 (20)	250 (13)	476 (29)	12 (14)	0*	0*	4.31
Br(7)	8f	0.0*	0.20916 (28)	0.36026 (46)	336 (25)	358 (17)	808 (44)	0*	0*	278 (22)	6.94
Br(8)	8f	0.0*	0.36035 (28)	0.35771 (47)	305 (26)	401 (19)	782 (46)	0*	0*	-318 (23)	7.05
Br(9)	8g	0.15928 (30)	0.24687 (23)	0.75*	203 (19)	233 (13)	430 (27)	-3 (15)	0*	0*	4.11
Br(10)	8f	0.0*	0.17461 (26)	0.87718 (46)	315 (26)	336 (16)	746 (42)	0*	0*	224 (21)	6.47
Br(11)	8f	0.0*	0.33738 (24)	0.86916 (41)	342 (27)	302 (15)	528 (34)	0*	0*	-128 (17)	5.55

Group	Position of centroid	$x_0^e$	$y_0$	$z_0$	$\delta$	$\epsilon$	$\eta$	Group $B,^b \text{ \AA}^2$
Ring 1	16h	0.2532 (12)	0.1674 (5)	0.4959 (14)	1.641 (18)	-2.609 (20)	-1.169 (15)	7.4 (4)
Ring 2	4a	0.0*	0.0*	0.5*	3.14159*	2.893 (38)	-1.57080*	7.9 (10)
Ring 3	4b	0.5*	0.0*	0.5*	-1.57080*	3.14159*	-1.768 (34)	9.7 (12)

<sup>a</sup> Numbers in parentheses in all tables and in the text are estimated standard deviations occurring in the least significant digit of the parameter. <sup>b</sup> The form of the anisotropic thermal ellipsoid is  $\exp[-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)]$ . The  $\beta$ 's and their estimated standard deviations have been multiplied by  $10^6$ . <sup>c</sup> Calculated from the anisotropic thermal parameters and unit cell parameters by the equation  $B \sim \frac{1}{3}(\beta_{11}a^2 + \beta_{22}b^2 + \beta_{33}c^2 + 2\beta_{12}ab \cos \gamma + 2\beta_{13}ac \cos \beta + 2\beta_{23}bc \cos \alpha)$ : W. C. Hamilton, *Acta Crystallogr.*, **12**, 609 (1959). <sup>d</sup> An asterisk denotes a parameter fixed by symmetry. <sup>e</sup>  $x_0, y_0, z_0$  are fractional coordinates of the ring centers. The angles  $\delta, \epsilon, \eta$  (in radians) are those described elsewhere: ref 27; R. Eisenberg and J. A. Ibers, *Inorg. Chem.*, **4**, 773 (1965). <sup>f</sup> Group  $B$  is the isotropic thermal parameter for the entire pyridinium cation.

 TABLE II  
 DERIVED POSITIONAL PARAMETERS FOR GROUP ATOMS IN  
 $(C_5H_5NH)_6Sb_4Br_{24}^a$ 

Atom	Position	$x$	$y$	$z$
Ring 1				
C(1)	16h	0.2143	0.1881	0.4294
C(2)	16h	0.2923	0.1808	0.4257
C(3)	16h	0.3312	0.1600	0.4923
C(4)	16h	0.2921	0.1466	0.5625
C(5)	16h	0.2140	0.1540	0.5662
C(6)	16h	0.1751	0.1747	0.4996
H(1)	16h	0.1866	0.2029	0.3820
H(2)	16h	0.3202	0.1903	0.3757
H(3)	16h	0.3868	0.1548	0.4897
H(4)	16h	0.3198	0.1318	0.6099
H(5)	16h	0.1861	0.1444	0.6162
H(6)	16h	0.1195	0.1799	0.5022
Ring 2				
C(7)	8f	0.0	0.0143	0.4186
C(8)	8f	0.0	0.0559	0.4772
C(9)	8f	0.0	0.0416	0.5586
H(7)	8f	0.0	0.0245	0.3607
H(8)	8f	0.0	0.0957	0.4610
H(9)	8f	0.0	0.0712	0.6003
Ring 3				
C(10)	8f	0.5	0.0114	0.4077
C(11)	16h	0.4318	0.0057	0.4588
H(10)	8f	0.5	0.0194	0.3590
H(11)	16h	0.3832	0.0098	0.4295

<sup>a</sup> Only positions of the crystallographically independent atoms are indicated, though entire rings were included in the refinement. Average estimated standard deviations ( $\times 10^4$ ) in parameters not fixed by symmetry: C,  $\sigma(x) = 16, \sigma(y) = 11, \sigma(z) = 19$ ; H,  $\sigma(x) = 22, \sigma(y) = 17, \sigma(z) = 25$ .

variables (95). On the final cycle the shift in each positional, thermal, and group parameter averaged 0.004 times its own  $\sigma$ . A final difference synthesis revealed no peaks greater than  $0.3 \text{ e}^-/\text{\AA}^3$ , consistent with good refinement. The largest peaks in this map were associated with the ring carbon atoms, positioned as to suggest slight rotational disorder of the cations about their ring normals.

The final positional, thermal, and group parameters derived from the last cycle of least-squares refinement are presented in Table I, along with the associated standard deviations in these parameters as estimated from the inverse matrix. The positional parameters of the ring carbon atoms, which may be derived from the data in Table I, are presented in Table II. Root-

 TABLE III  
 FINAL ROOT-MEAN-SQUARE THERMAL AMPLITUDES OF VIBRATION  
 ( $\text{\AA}$ ) IN  $(C_5H_5NH)_6Sb_4Br_{24}$ 

Atom	Min	Med	Max
Sb(1)	0.183 (4)	0.211 (4)	0.216 (5)
Sb(2)	0.166 (9)	0.206 (7)	0.230 (8)
Sb(3)	0.166 (8)	0.191 (7)	0.202 (7)
Br(1)	0.208 (10)	0.279 (9)	0.319 (8)
Br(2)	0.224 (9)	0.281 (9)	0.296 (8)
Br(3)	0.203 (9)	0.276 (8)	0.284 (8)
Br(4)	0.222 (9)	0.277 (9)	0.319 (9)
Br(5)	0.186 (5)	0.283 (5)	0.302 (7)
Br(6)	0.166 (9)	0.254 (8)	0.267 (7)
Br(7)	0.226 (8)	0.228 (9)	0.401 (9)
Br(8)	0.217 (9)	0.218 (8)	0.416 (9)
Br(9)	0.177 (8)	0.242 (8)	0.258 (7)
Br(10)	0.221 (9)	0.233 (8)	0.378 (8)
Br(11)	0.230 (8)	0.230 (9)	0.324 (8)



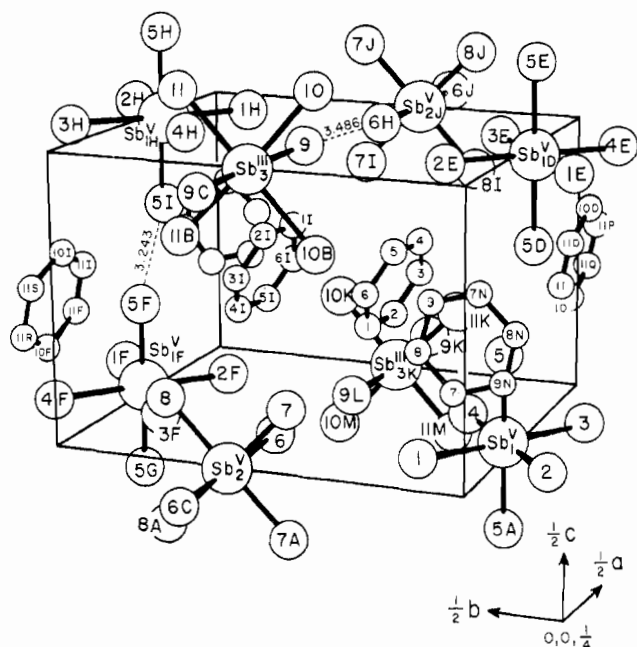


Figure 2.—Crystal structure of  $(C_5H_5NH)_6Sb_4Br_{24}$ , illustrating the contents of one-eighth of the unit cell (two asymmetric units). The view is identical with that in Figure 1. For clarity, pyridinium hydrogen atoms have been excluded. The numbering scheme of all atoms, including those related by symmetry, coincides with that in Tables I, II, and V (footnote *b*).

TABLE V  
DIMENSIONS OF THE ANIONS IN  $(C_5H_5NH)_6Sb_4Br_{24}$

Atoms	Length, Å		Atoms <sup>b</sup>	Angle, deg
	Uncor	Cor <sup>a</sup>		
(a) $Sb^VBr_6^-$ Anion (Position 8g)				
Sb(1)—Br(1)	2.531 (7)	2.555 (10)	Br(1)—Sb(1)—Br(2)	92.2 (2)
Sb(1)—Br(2)	2.532 (7)	2.547 (10)	Br(1)—Sb(1)—Br(4)	90.1 (2)
Sb(1)—Br(3)	2.566 (7)	2.580 (10)	Br(1)—Sb(1)—Br(5)	90.0 (1)
Sb(1)—Br(4)	2.522 (7)	2.537 (10)	Br(2)—Sb(1)—Br(3)	89.7 (2)
Sb(1)—Br(5)	2.532 (4)	2.544 (6)	Br(2)—Sb(1)—Br(5)	90.1 (1)
			Br(3)—Sb(1)—Br(4)	88.0 (2)
			Br(3)—Sb(1)—Br(5)	90.0 (1)
			Br(4)—Sb(1)—Br(5)	89.0 (1)
(b) $Sb^VBr_6^-$ Anion (Position 4c)				
Sb(2)—Br(6)	2.535 (5)	2.542 (7)	Br(6)—Sb(2)—Br(7)	89.2 (1)
Sb(2)—Br(7)	2.531 (7)	2.546 (10)	Br(6)—Sb(2)—Br(8)	90.9 (1)
Sb(2)—Br(8)	2.529 (7)	2.545 (10)	Br(7)—Sb(2)—Br(7A)	91.1 (4)
			Br(7)—Sb(2)—Br(8)	90.2 (3)
			Br(8)—Sb(2)—Br(8A)	88.5 (4)
(c) $Sb^{III}Br_6^{3-}$ Anion (Position 4c)				
Sb(3)—Br(9)	2.788 (5)	2.793 (7)	Br(9)—Sb(3)—Br(10)	87.9 (1)
Sb(3)—Br(10)	2.797 (7)	2.801 (10)	Br(9)—Sb(3)—Br(11)	92.2 (1)
Sb(3)—Br(11)	2.789 (7)	2.809 (10)	Br(10)—Sb(3)—Br(10B)	96.3 (3)
			Br(10)—Sb(3)—Br(11)	87.4 (2)
			Br(11)—Sb(3)—Br(11B)	88.9 (3)

<sup>a</sup> The estimated standard deviations include an arbitrary uncertainty factor of 1.4 for the rigid-body assumption. <sup>b</sup> Key to symmetry operations in this table, succeeding tables, Figure 2, and the text: (A)  $x, y, 1/2 - z$ ; (B)  $x, y, 1.5 - z$ ; (C)  $\bar{x}, y, z$ ; (D)  $x, \bar{y}, 1 - z$ ; (E)  $x, \bar{y}, 1/2 + z$ ; (F)  $1/2 - x, 1/2 + y, z$ ; (G)  $1/2 - x, 1/2 + y, 1/2 - z$ ; (H)  $1/2 - x, 1/2 - y, 1/2 + z$ ; (I)  $1/2 - x, 1/2 - y, 1 - z$ ; (J)  $1/2 + x, 1/2 - y, 1/2 + z$ ; (K)  $1/2 + x, 1/2 - y, -1/2 + z$ ; (L)  $1/2 - x, 1/2 - y, -1/2 + z$ ; (M)  $1/2 + x, 1/2 - y, 1 - z$ ; (N)  $\bar{x}, \bar{y}, 1 - z$ ; (P)  $1 - x, \bar{y}, 1 - z$ ; (Q)  $1 - x, y, z$ ; (R)  $-1/2 + x, 1/2 + y, z$ ; (S)  $-1/2 + x, 1/2 - y, 1 - z$ .

pyridinium salt are essentially equivalent within one standard deviation, and deviations of the bond angles from  $90^\circ$  do not exceed  $2.2^\circ$ . One bond, 2.580 (10) Å in length, in the  $Sb^VBr_6^-$  ion of Sb(1) and one angle,

TABLE VI  
AVERAGE Sb—Br BOND LENGTHS (Å) IN  
 $Sb^VBr_6^-$  AND  $Sb^{III}Br_6^{3-}$  IONS<sup>a</sup>

Compound	$Sb^VBr_6^-$	$Sb^{III}Br_6^{3-}$	Ref
$(NH_4)_6Sb_4Br_{24}$ <sup>b</sup>	2.552 (10)	2.792 (4)	<i>d</i>
$Rb_4Sb_2Br_{12}$ <sup>c</sup>	2.550 (47)	2.763 (37)	<i>e</i>
$(C_5H_5NH)_6Sb_4Br_{24}$	2.547 (16), 2.544 (2)	2.799 (7)	This study
$(\alpha-C_6H_7NH)_2SbBr_9$	2.553 (9)	...	<i>f</i>
Av	2.55 (2)	2.78 (3)	

<sup>a</sup> Bond lengths in the ions in all four structures have been corrected for thermal effects using the Cruickshank approximation of rigid-body libration.<sup>23</sup> The breadth parameter,  $q^2$ , was assigned a value of  $0.12 \text{ \AA}^2$  in each case. Average distances and associated rms deviations were computed from the expressions

$$\bar{x} = \frac{\sum_{i=1}^N (x_i/\sigma_i^2)}{\sum_{i=1}^N (1/\sigma_i^2)}$$

$$\sigma(\bar{x}) = \left[ \frac{\sum_{i=1}^N (x_i - \bar{x})^2 / (N - 1)}{\sum_{i=1}^N (1/\sigma_i^2)} \right]^{1/2}$$

where  $x_i$  is an individual observation,  $\sigma_i$  is the corresponding standard deviation, and  $N$  is the number of observations. Averages for a given ion include the lengths of all six bonds. <sup>b</sup> The corrected bond lengths previously reported<sup>5-7</sup> were based on the Busing-Levy "riding" model, not the rigid-body librational model as had been reported. The corrected bond lengths of the independent bonds based on the rigid-body assumption are Sb(V)—Br(1) = 2.560 (8), Sb(V)—Br(3) = 2.541 (7), Sb(III)—Br(2) = 2.789 (8), and Sb(III)—Br(4) = 2.796 (7) Å. <sup>c</sup> The corrected bond lengths of the independent bonds based on the rigid-body assumption are Sb(V)—Br(1) = 2.536 (7), Sb(V)—Br(3) = 2.622 (11), Sb(III)—Br(2) = 2.787 (6), and Sb(III)—Br(4) = 2.716 (6) Å. <sup>d</sup> See ref 6; also footnote *b*, this table. <sup>e</sup> See ref 7; also footnote *c*, this table. <sup>f</sup> See ref 8.

TABLE VII  
SELECTED NEAREST-NEIGHBOR C...Br, H...Br, AND Br...Br CONTACTS AND ANGLES

Atoms	Dist. <sup>a</sup> Å	Atoms	Dist. <sup>a</sup> Å
Carbon...Bromine Contacts			
C(1)...Br(1)	4.09	C(5)...Br(9)	3.85
C(1)...Br(4)	4.14	C(6)...Br(5)	4.11
C(1)...Br(5)	4.06	C(7)...Br(1)	3.90
C(2)...Br(4)	3.57	C(7)...Br(2)	3.95
C(2)...Br(5)	4.07	C(7)...Br(5)	3.87
C(3)...Br(5)	4.14	C(8)...Br(5)	4.14
C(4)...Br(5)	4.19	C(10D)...Br(8I)	3.71
C(4)...Br(5D)	4.12	C(10)...Br(11K)	3.67
C(4)...Br(6H)	3.68	C(11D)...Br(8I)	4.00
C(5)...Br(5)	4.18	C(11)...Br(11K)	4.17
C(5)...Br(5D)	4.10		
Hydrogen...Bromine Contacts			
H(1)...Br(6)	2.86	H(6)...Br(10B)	2.80
H(2)...Br(4)	3.03	H(7)...Br(1)	3.10
H(2)...Br(9L)	2.48	H(8)...Br(7)	3.06
H(3)...Br(11K)	2.73	H(9)...Br(10B)	2.40
H(4)...Br(3E)	2.85	H(11)...Br(5)	2.77
H(5)...Br(2E)	3.03		
Bromine...Bromine Contacts			
Br(1)...Br(1C)	4.083 (14)	Br(3F)...Br(8)	3.998 (8)
Br(1)...Br(6)	4.524 (8)	Br(4)...Br(9L)	3.781 (9)
Br(1)...Br(7)	3.893 (8)	Br(4)...Br(11K)	4.022 (7)
Br(2)...Br(2C)	4.348 (14)	Br(5)...Br(5D)	3.243 (7)
Br(2E)...Br(9)	4.424 (8)	Br(6)...Br(9L)	3.486 (7)
Br(2E)...Br(10)	4.025 (8)	Br(7)...Br(10B)	4.381 (10)
Br(3F)...Br(6)	3.898 (8)	Br(8)...Br(11B)	4.510 (10)
Angles, deg			
Br(5)—Sb(1)—Br(5A)	179.8 (3)	Sb(1)—Br(5)—Br(5D)	164.8 (3)
Br(6)—Sb(2)—Br(6C)	177.6 (3)	Sb(2)—Br(6)—Br(9L)	176.9 (3)
Br(9)—Sb(3)—Br(9C)	173.7 (2)	Sb(3)—Br(9)—Br(8H)	171.9 (3)

<sup>a</sup> Distances and angles do not include a correction for thermal effects of vibrating anions. The H...Br distances are based on C—H bond lengths of 1.1 Å. Average estimated standard deviations in the C...Br distances are 0.03 Å and those in the H...Br distances are 0.06 Å.

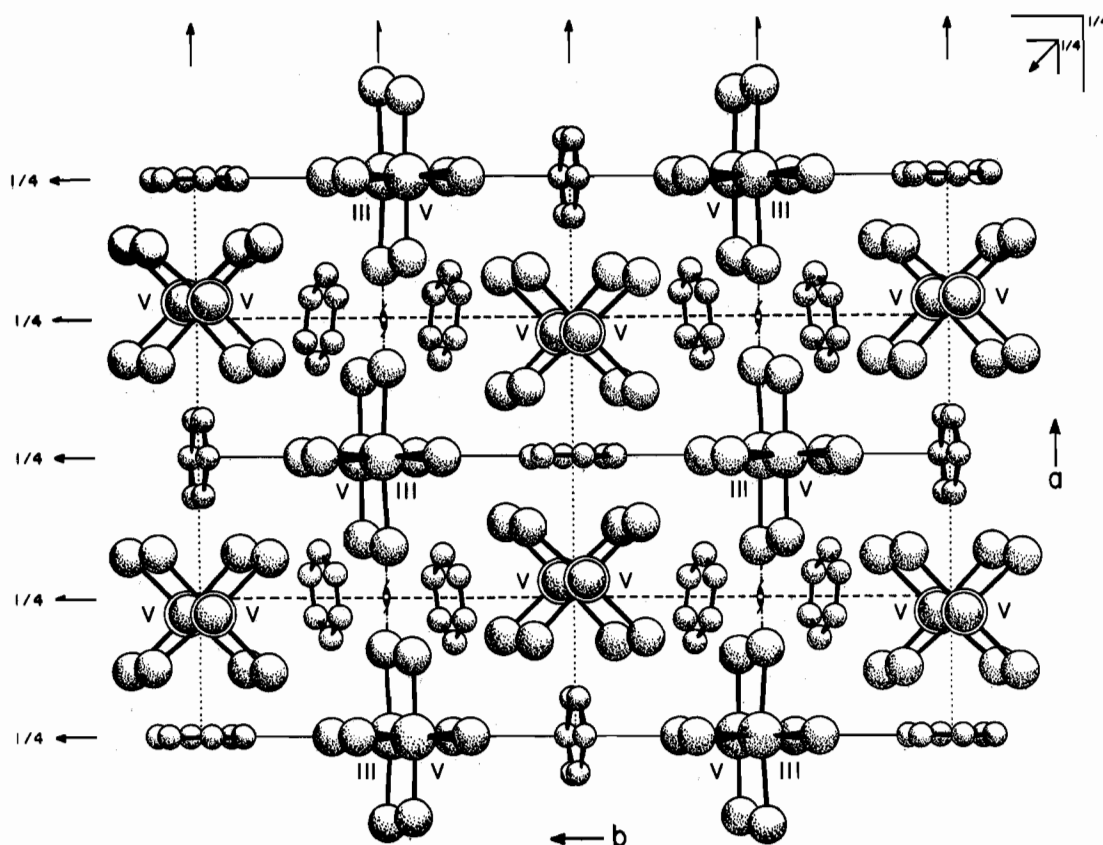


Figure 3.—Crystal structure of  $(C_5H_5NH)_6Sb_4Br_{24}$  projected onto the (001) plane. Only the anions in levels  $z = 1/4$  and  $3/4$  and the cations in level  $z = 1/2$  within the unit cell are shown. For clarity, pyridinium hydrogen atoms have been excluded.

96.3 (3)°, in the  $Sb^{III}Br_6^{3-}$  ion are the only two exceptions.

Selected nearest-neighbor carbon···bromine and hydrogen···bromine distances are given in Table VII. By comparison with the 3.65-Å<sup>34</sup> sum of the Pauling van der Waals radius of bromine and the half-thickness of the aromatic pyridinium ring, it is evident that no unusual cation···bromine distances along the ring normals are present. The closest approach of any ring is 3.87 Å, the distance of Br(5) from the plane of ring II. Eleven cationic hydrogen···bromine distances less than the Pauling van der Waals radius sums of hydrogen and bromine, 3.2 Å,<sup>34</sup> are present, however. The shortest distance, 2.4 Å, involves H(9) with Br(10B).

Nearest-neighbor interionic bromine···bromine contacts are also listed in Table VII. Two distances significantly shorter than the 3.90-Å van der Waals radius sums are present in the structure. These distances are 3.486 (7) Å between Sb(V) and Sb(III) anions (chain I) and 3.243 (7) Å between neighboring Sb(V) anions (chain II). These distances strongly indicate the presence of weak bonding interactions between these ions. Such short distances, particularly the 3.24-Å distance, would be energetically unfavorable in the absence of some bond formation, indicated as a first approximation by calculation of the interaction energy using the Lennard-Jones potential. Attractive interactions of this type are already known. In solid bromine, for example, intermolecular distances as short as

(34) L. Pauling, "The Nature of the Chemical Bond," Cornell University Press, Ithaca, N. Y., 1960, p 260.

3.30 Å occur.<sup>35</sup> All Br-Br···Br valency angles are close to 90° or to 180°, as in our case, a structural feature found among all polyhalide complexes except  $IF_7$ .<sup>36-38</sup> An even shorter distance, 2.87 Å, has been found between  $Br_2$  molecules and  $Sb_2Br_9^{3-}$  ions in  $[N(CH_3)_4]_8-Sb_2Br_{11}$ .<sup>39</sup> Also noteworthy is that the interaction along chain I is virtually as short as the Br···Br distance, 3.491 (2) Å, found within the nearly linear  $\cdots Br_3 \cdots Sb^V Br_6 \cdots$  chain in  $(\alpha-C_6H_7NH)_2SbBr_9$ . This remarkable similarity of the two interactions, summarized in Figure 4, is readily appreciated by considering the structural similarities of the tribromide and  $Sb^{III}-Br_6^{3-}$  ions. First, both ions exhibit similar bonding; expressed in terms of the LCAO-MO treatment, the molecular orbitals may be represented as linear combinations of the outermost p orbitals of the bromine atoms in the normal three-center, four-electron approximation.<sup>9,36,37,40</sup> Second, the external bromines in both ions carry a formal charge of  $1/2-$ .

Aromatic molecules are known to form donor-acceptor adducts with such molecules as bromine<sup>41</sup> and carbon tetrabromide.<sup>42</sup> In both cases the distance of bromine from the ring centers is approximately 0.3 Å shorter

(35) B. Vonnegut and B. E. Warren, *J. Amer. Chem. Soc.*, **55**, 2459 (1936).

(36) E. E. Havinga and E. H. Wiebenga, *Recl. Trav. Chim. Pays-Bas*, **78**, 724 (1959).

(37) E. H. Wiebenga, E. E. Havinga, and K. H. Boswijk, *Advan. Inorg. Chem. Radiochem.*, **3**, 133 (1961).

(38) H. A. Bent, *Chem. Rev.*, **63**, 587 (1963).

(39) C. R. Hubbard and R. A. Jacobson, to be submitted for publication.

(40) R. E. Rundle, *J. Amer. Chem. Soc.*, **85**, 112 (1963).

(41) O. Hassel and K. O. Strømme, *Acta Chem. Scand.*, **12**, 1146 (1958).

(42) F. J. Strieter and D. H. Templeton, *J. Chem. Phys.*, **37**, 161 (1962).

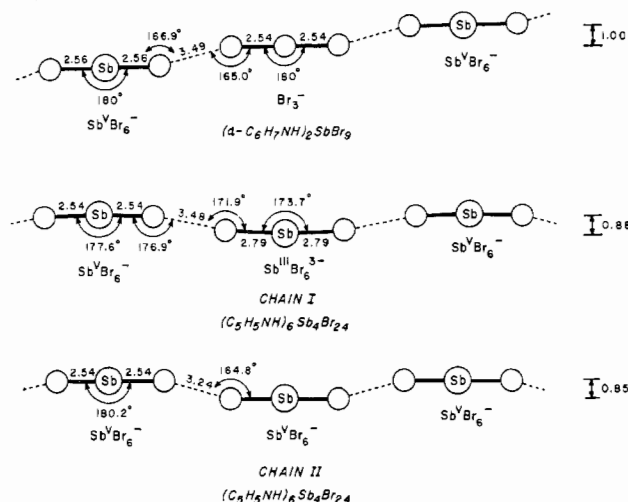


Figure 4.—Schematic representation of the nearly linear anionic chains in  $(\alpha\text{-C}_6\text{H}_7\text{NH})_2\text{SbBr}_9$  and  $(\text{C}_5\text{H}_5\text{NH})_6\text{Sb}_4\text{Br}_{24}$ . Interatomic distances, angles, and the perpendicular displacements of the anionic centers along the chains are indicated.

than the conventional van der Waals distance (3.65 Å). In our material, three cations surrounding the  $\text{Br}\cdots\text{Br}$  bond in chain II face bromine atoms. However, all distances exceed van der Waals distances, leading us to conclude that donor-acceptor interactions between the pyridinium cations and the  $\text{Sb}^{\text{V}}\text{Br}_6^-$  anions are very weak or, more likely, nonexistent.

Unlike the deep red 2-, 3-, and 4-methylpyridinium salts which undergo color changes at the temperature of liquid nitrogen,<sup>11</sup> the deep red color of this material is unaffected by a decrease in temperature. This additional absorption by the pyridinium derivative may be ascribed to possible charge-transfer absorption within the anions themselves or to possible intervalence-transfer absorption between the unlike anions.<sup>43</sup> Since the absorption properties of the individual anions are unknown, and in view of the weak bonding between the mixed Sb(III) and Sb(V) moieties, bonding not found in the ammonium and rubidium analogs, the relative importance of intervalence-transfer absorption in this material cannot be fully assessed. Further research in this area is certainly dictated.

Deformation of the  $\text{Sb}^{\text{III}}\text{Br}_6^{3-}$  ion from  $O_h$  symmetry immediately invites speculation regarding the stereochemical "activity" of a lone pair of electrons. This hexahalide ion is of considerable interest because the valence shell of the central antimony atom contains seven electron pairs of which only six are actually used in bonding. According to the simple valence-shell

(43) Three extensive reviews of mixed-valence chemistry, including the Sb(III)-Sb(V) system, have recently been published. They survey in considerable detail all compounds presently known to exhibit intervalence electron-transfer absorption. Theoretical aspects of this phenomenon are also presented: (a) M. B. Robin and P. Day, *Advan. Inorg. Chem. Radiochem.*, **10**, 247 (1967); (b) G. C. Allen and N. S. Hush, *Progr. Inorg. Chem.*, **8**, 357 (1967); (c) N. S. Hush, *ibid.*, **8**, 392 (1967). Detailed studies of the electronic spectra and semiconductivity of hexachloroantimonates(III,V) have also been reported: (d) L. Atkinson and P. Day, *J. Chem. Soc. A*, 2423 (1969); (e) L. Atkinson and P. Day, *ibid.*, 2432 (1969).

electron-pair repulsion model of Sidgwick and Powell<sup>44</sup> and Gillespie and Nyholm,<sup>45</sup> lone pairs as well as bonded electron pairs influence the stereochemistry of molecules and ions. It has been shown that for the seven-coordination case more than one arrangement of electron pairs with similar energies is possible.<sup>45b,46-48</sup> Two of these arrangements are exhibited by  $\text{XeF}_6$ <sup>49</sup> and  $\text{IF}_7$ .<sup>50-52</sup> The molecule  $\text{XeF}_6$ , which is formally isoelectronic with the  $\text{Sb}^{\text{III}}\text{Br}_6^{3-}$  ion and contains a lone pair of electrons, exists as an irregular octahedron exhibiting a preference toward  $C_{3v}$  symmetry, though it is considerably less distorted than predicted solely by the electron-pair repulsion model. The hexahalide ions  $\text{SeCl}_6^{2-}$ ,  $\text{SeBr}_6^{2-}$ ,  $\text{TeCl}_6^{2-}$ , and  $\text{TeBr}_6^{2-}$ , on the other hand, each exhibit *regular* octahedral structures and therefore constitute definite exceptions to Gillespie's hypothesis,<sup>53-58</sup> similarly, the  $\text{SbBr}_6^{3-}$  ion in  $(\text{NH}_4)_4\text{-Sb}_2\text{Br}_{12}^6$  and  $\text{Rb}_4\text{Sb}_2\text{Br}_{12}^7$  also exhibits no appreciable deformation. This lack of deformation in the case of the ions may be attributed to ligand-ligand repulsion arising from steric effects of the relatively large chlorine and bromine atoms, forces which are sufficiently strong to prevent virtually any stereochemical influence of the lone pair.<sup>45b</sup> The most prominent deformation of the ion corresponds to a slight opening of a  $\text{Br-Sb-Br}$  angle to  $96.3(3)^\circ$ , but there is no associated compression of the remaining four bromine ligands toward the opposite side of the ion. We therefore ascribe the increase in the angle  $\text{Br}(10)\text{-Sb}(3)\text{-Br}(10\text{B})$  from  $90$  to  $96.3^\circ$  primarily to ligand-ligand repulsion introduced by the decrease in the angle  $\text{Br}(9)\text{-Sb}(3)\text{-Br}(9\text{C})$  from  $180$  to  $173.7(3)^\circ$ , a decrease presumably caused by the weak bond formation between the ions along chain I. Possible weak hydrogen  $\cdots$  bromine bonding between  $\text{Br}(10\text{B})$  and  $\text{H}(9)$  at a distance of  $2.4$  Å may also contribute to the ion deformation.

**Acknowledgments.**—Helpful discussions with Drs. L. S. Bartell and D. E. Williams are gratefully acknowledged.

(44) N. V. Sidgwick and H. M. Powell, *Proc. Roy. Soc., Ser. A*, **176**, 153 (1940).

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(46) D. Britton, *Can. J. Chem.*, **41**, 1632 (1963).

(47) (a) T. A. Claxton and G. C. Benson, *ibid.*, **44**, 157 (1966); (b) T. A. Claxton and G. C. Benson, *ibid.*, **44**, 1730 (1966).

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